

Recent results in Color Superconductivity

G. Nardulli

QCD Phase at $T=0$, large μ

Gluonic interaction between two quarks

$$3 \times 3 = \bar{3} + 6$$

The color antisymmetric $\bar{3}$ channel is attractive. Therefore Cooper's theorem implies formation of a condensate of Cooper pairs of quarks (diquarks) and color superconductivity

$$\langle \psi\psi \rangle = \Delta \neq 0$$

Color Superconductivity at extreme densities

- **S=0** state is preferred as more phase space available
- Pauli principle implies antisymmetry in flavor
- Asymptotically (very large densities): Color Flavor Locking (CFL; Alford, Rajagopal, Wilczek)

$$\langle 0 | \psi^i_{\alpha} \psi^j_{\beta} | 0 \rangle = \Delta \epsilon^{ij\gamma} \epsilon_{\alpha\beta\gamma}$$

Valid for $\mu \gg m_q$

Rigourously proved in QCD, albeit at very large μ .

At intermediate densities, relevant for astrophysical applications the ground state is still unknown

Model calculations (NJL) as QCD unreliable

Deviations from CFL due to s quark mass effect.

**Color and electric neutrality must be imposed;
=> different gaps for quarks of different flavor.**

Gapless phases uniform (gCFL) or not uniform (LOFF).

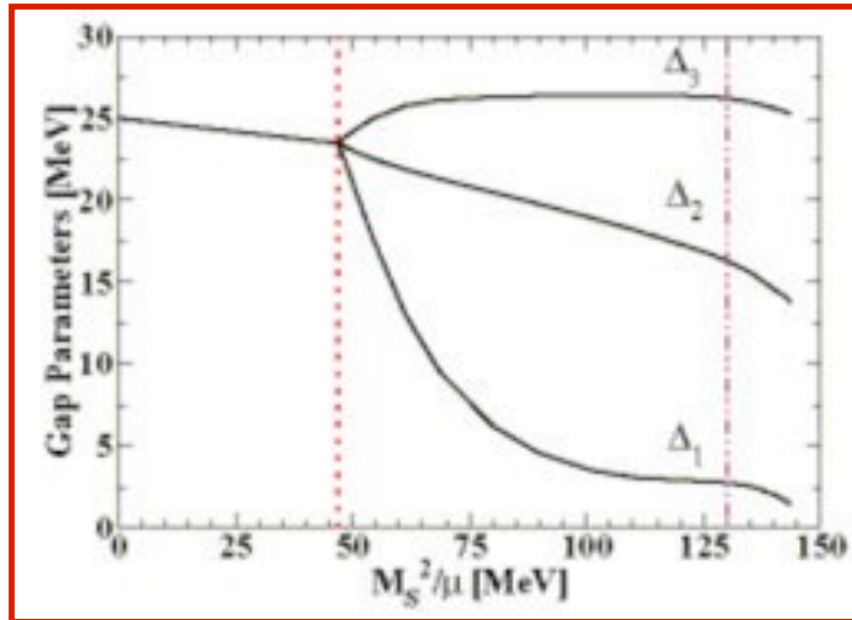
Choosing the true vacuum

The true ground state has the minimum value of the free energy (grand potential)

$$\Omega = -\frac{1}{2\pi^2} \int_0^\Lambda dp p^2 \sum_{j=1}^9 |\epsilon_j(\ell_{\parallel})| + \frac{1}{G} (\Delta_1^2 + \Delta_2^2 + \Delta_3^2) - \frac{\mu_Q^4}{12\pi^2}$$

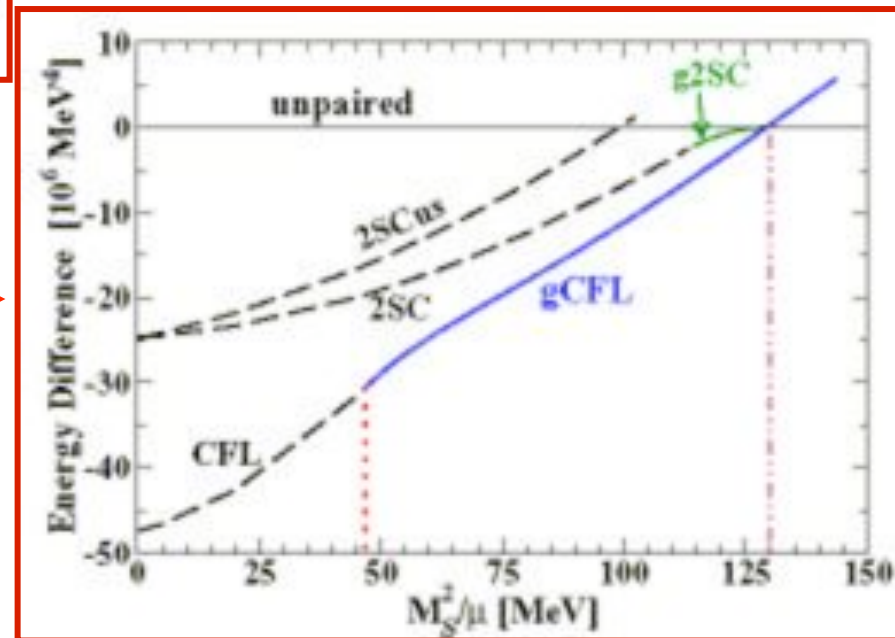
Results for gCFL phase

Gap Parameters



Free energy

M.Alford, P.Jotwani, C.
Kouvaris, J. Kundu,
K.Rajagopal

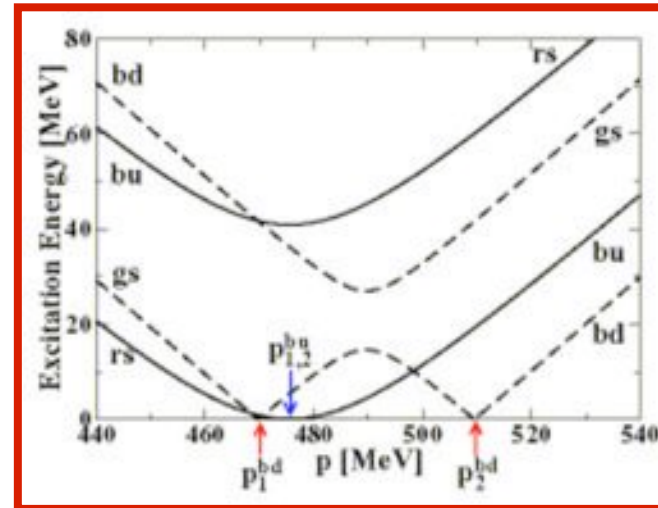


The importance of being gapless

$$\begin{aligned}\mu &= 500 \text{ MeV} \\ \Delta &= 25 \text{ MeV} \\ M_s^2/2\mu &= 80 \text{ MeV}\end{aligned}$$

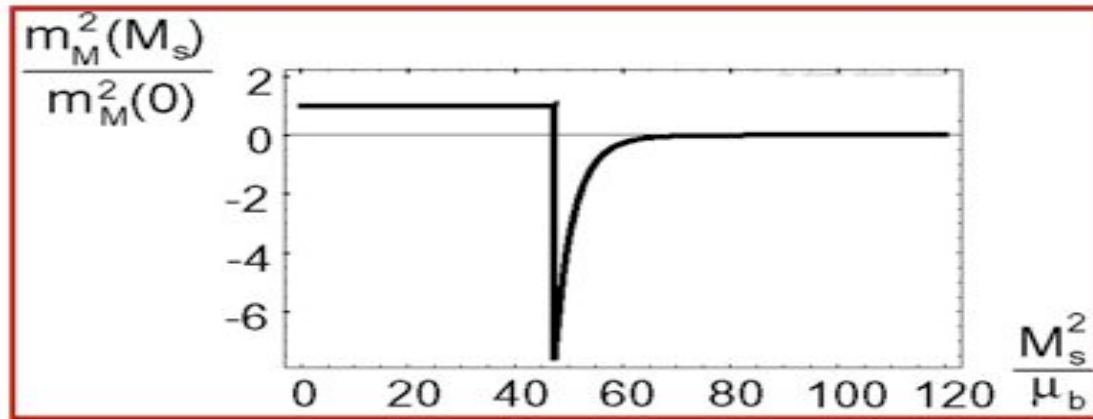
$$|\epsilon(p)| = |\delta\mu \pm \sqrt{(p - \mu)^2 + \Delta^2}|$$

$$p_{\text{gapless}} = \mu \pm \sqrt{\delta\mu^2 - \Delta^2}$$



**Gapless modes: relevant for astrophysical effects
(phenomenology of compact stars)**

Meissner masses in the gCFL phase



Solid line $a=1,2$;
One Gluon "3mixed8":
same behavior;
others: no instability
One $U_{em}(1)$: massless

R. Casalbuoni,
M. Mannarelli,
G. N., M. Ruggieri, R. Gatto

Imaginary masses: instability of chromomagnetic type
gCFL is not the true vacuum

Similar results for 5 gluons of broken generators in 2-flavor
(Huang-Shovkovy)

LOFF phase favored (Giannakis-Ren, Fukushima)

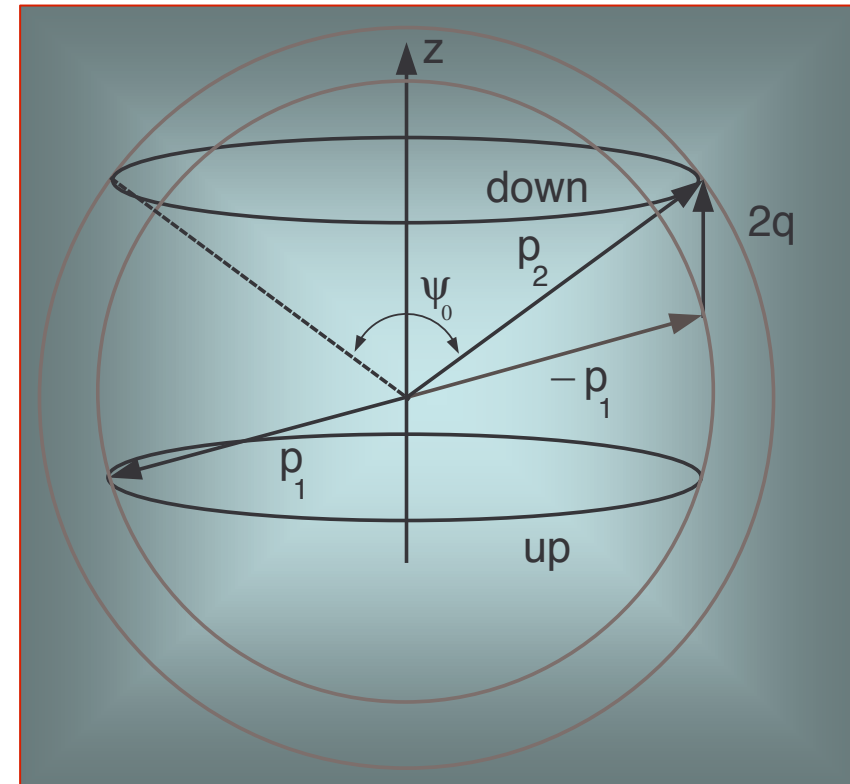
Loff Phase: two flavors (u,d)

If $\delta\mu \neq 0$ it can be energetically favourable to have states where the Cooper pair has total momentum $2q \neq 0$

LOFF= Larkin-Ovchinnikov & Fulde-Ferrell (1964);
Alford, Bowers, Rajagopal (2000)
R. Casalbuoni, G.N.
Rev. Mod. Phys.(2004)

Studied in condensed matter (2D superconductors), cuprates

$$\Delta(\mathbf{r}) = \langle 0 | \psi(\mathbf{r})\psi(\mathbf{r}) | 0 \rangle = \Delta \exp[i(\mathbf{p}_1 + \mathbf{p}_2) \cdot \mathbf{r}]$$



LOFF phase: three flavors

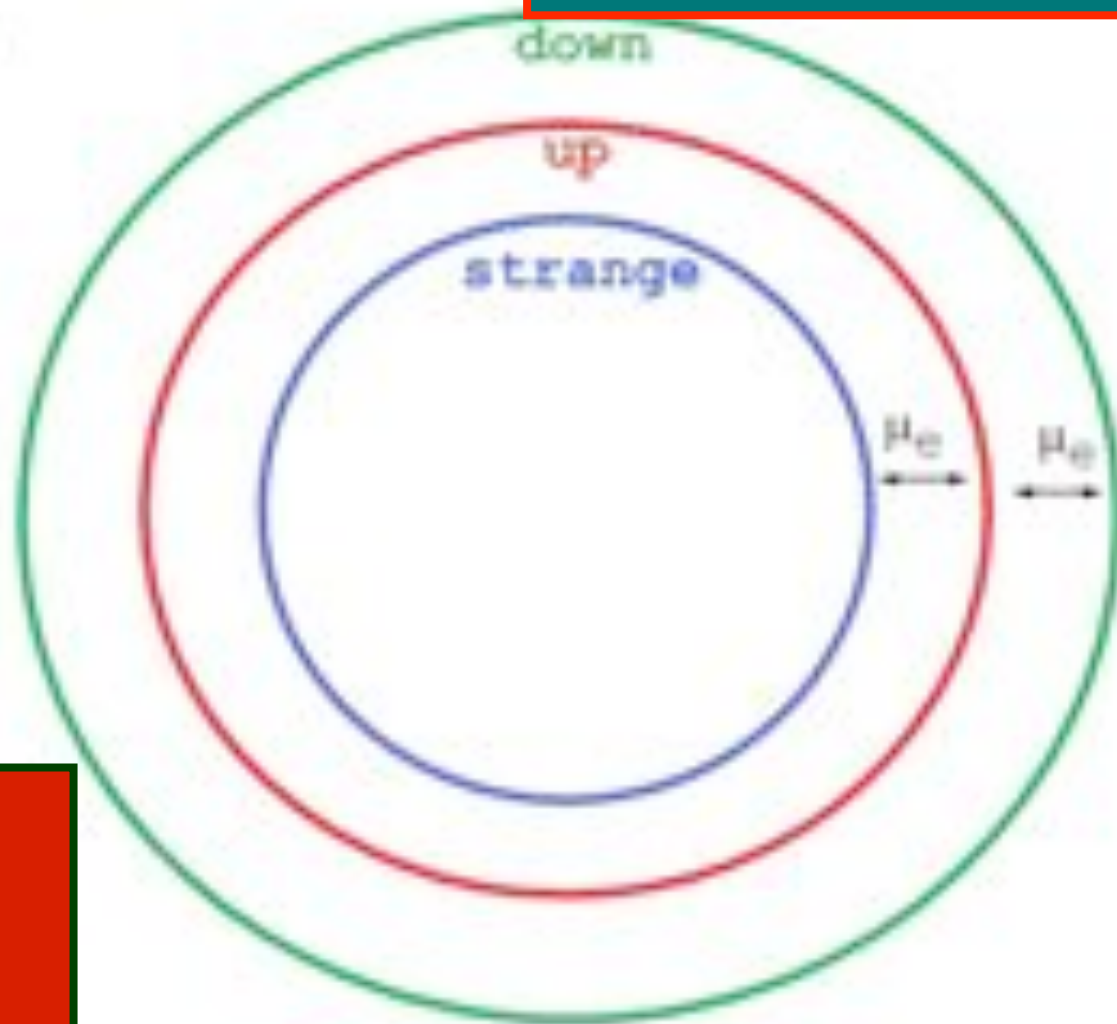
Needed for realistic calculations

- **R. Casalbuoni, R.Gatto,,N.Ippolito,GN, M.Ruggieri :
Ginzburg Landau expansion of the gap equation
and free energy**
- **(Almost) complete analysis beyond GL:
M.Mannarelli, K. Rajagopal, R. Sharma**

$$\mu_e = \mu_d - \mu_u = \mu_u - \mu_s \sim m_s^2 / 4\mu$$

$$\mu_d - \mu_s = -2\mu_e$$

Fermisurfaces
for quark
matter in the
normal phase,
 β equilibrium
and non zero
strange quark
mass



Near the transition point
 μ_u, μ_s couplings, equal
strength, no d_s pairing

Expected:
 $\mu_e = m_s^2 / 4\mu$
 $\Delta_1 = 0; \Delta_2 = \Delta_3$

Gap parameters

$$\Delta_{\mathbf{k}}(\mathbf{r}) = \Delta_{\mathbf{k}} e^{2i \mathbf{q}_{\mathbf{k}} \cdot \mathbf{r}}$$

For each inhomogeneous pairing a Fulde-Ferrell ansatz;

$$2 \mathbf{q}_{\mathbf{k}}$$

represents the momentum of the Cooper pair.

This is the simplest ansatz, other structures should be examined

Three independent functions

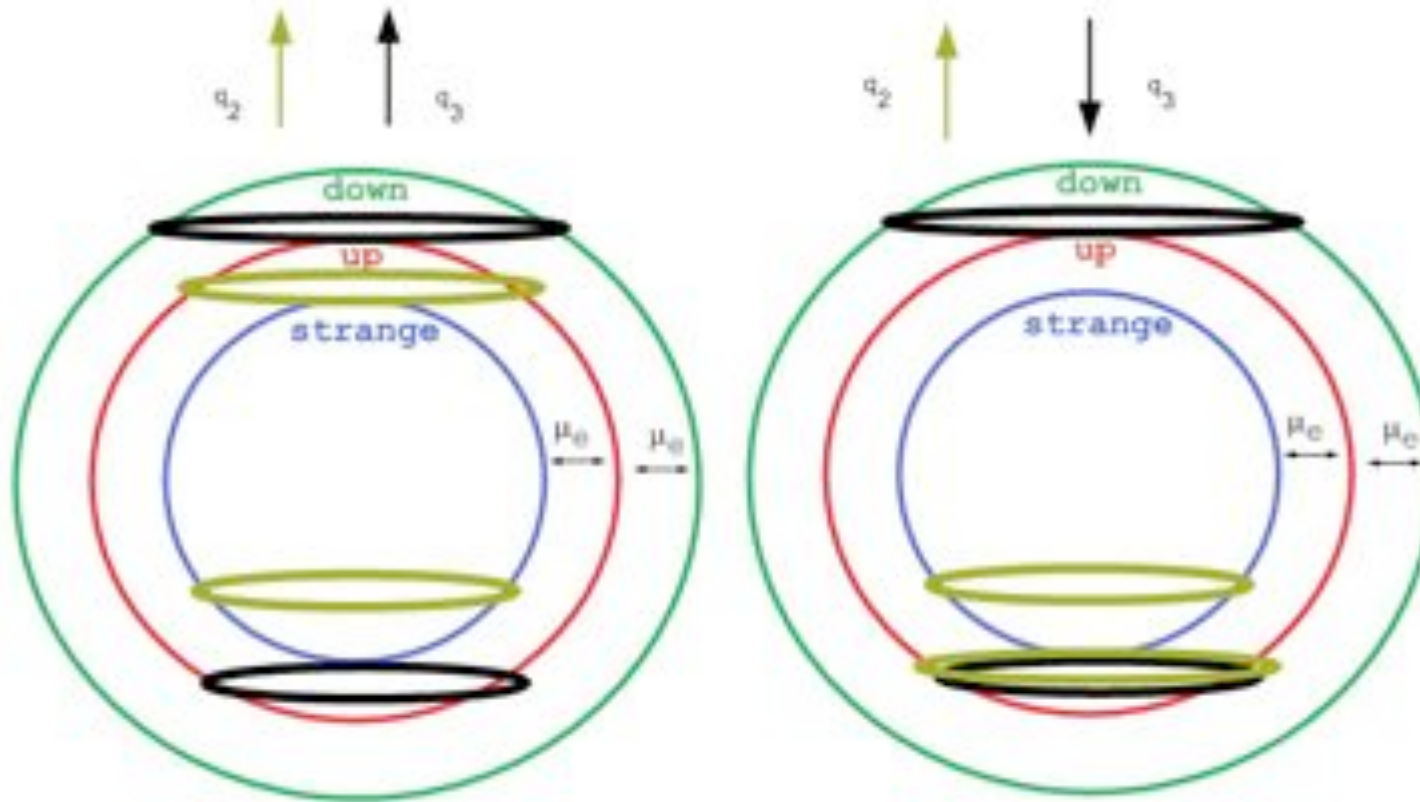
$$\Delta_1(\mathbf{r}),$$

$$\Delta_2(\mathbf{r}),$$

$$\Delta_3(\mathbf{r}),$$

describing respectively **d-s**, **u-s** and **u-d** pairing.

q_2 q_3 parallel favored by phase space



$$q \approx 1.2 \delta\mu$$

Ginzburg-Landau expansion

$$\Omega = \Omega_n + \frac{1}{4} \sum_{I=1}^3 (2\alpha_I \Delta_I^2 + \beta_I \Delta_I^4 + \sum_{J \neq I} \beta_{IJ} \Delta_I^2 \Delta_J^2) + O(\Delta^6)$$

$$\Omega_n = -\frac{3}{12\pi^2} (\mu_u^4 + \mu_d^4 + \mu_s^4) - \frac{1}{12\pi^2} \mu_e^4$$

$$-\partial_{\Delta_k} \Omega = 0 \text{ minimum in } \Delta_k$$

$$-\partial_{\mu_e} \Omega = 0 \text{ electrical neutrality}$$

Approximation: $\mu_3 \simeq \mu_8 \simeq 0$ (confirmed by R.Ciminale et al., in preparation: Valid near the transition point, since this is the result for the normal phase)

The numerical results in this regime confirm $\Delta_1 = 0$; $\Delta_2 \simeq \Delta_3$

R. Casalbuoni, R.Gatto,
N.Ippolito,GN,
M.Ruggieri

(Almost) complete calculation

Assume from the very beginning

$$\Delta_1 = 0 ; \quad \Delta_2 = \Delta_3$$
$$\mu_3 = \mu_8 = 0 ; \quad \mu_e = m_s^2 / 4\mu$$

Since

$$k_u + k_d = 2q_3 ; \quad k_u + k_s = 2q_2$$

Shift momenta so that

$$P_u = p = \mu v ; \quad P_d = p + 2q_3 ; \quad P_s = p + 2q_2$$

Corresponding to redefinitions of $\delta\mu$:

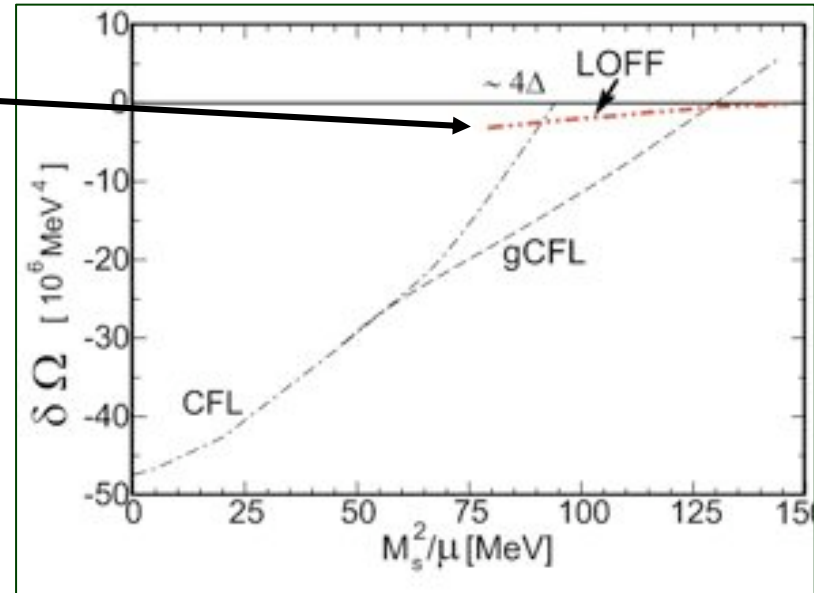
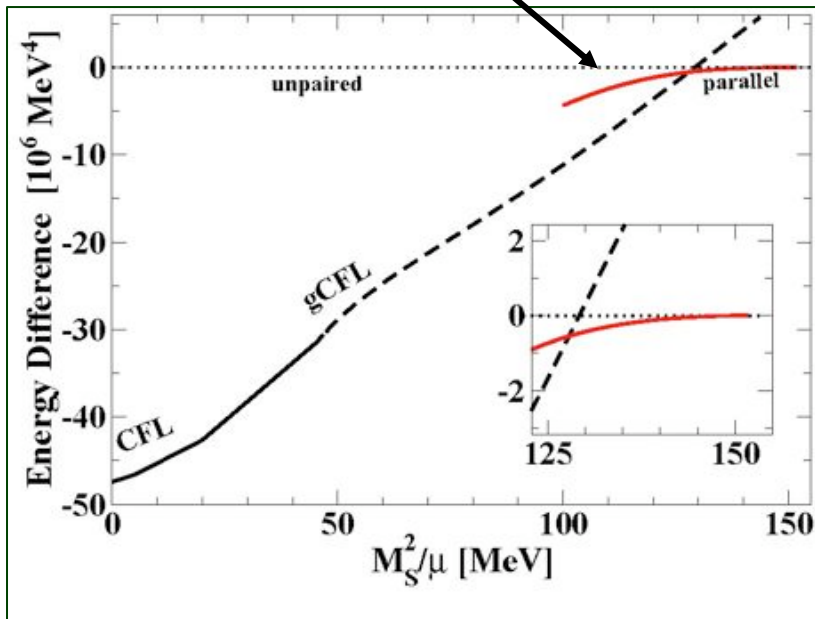
$$\delta\mu \rightarrow \delta\mu \pm q v$$

M. Mannarelli,
K. Rajagopal,
R. Sharma

Comparison among different CSC phases

Ginzburg Landau Approximation

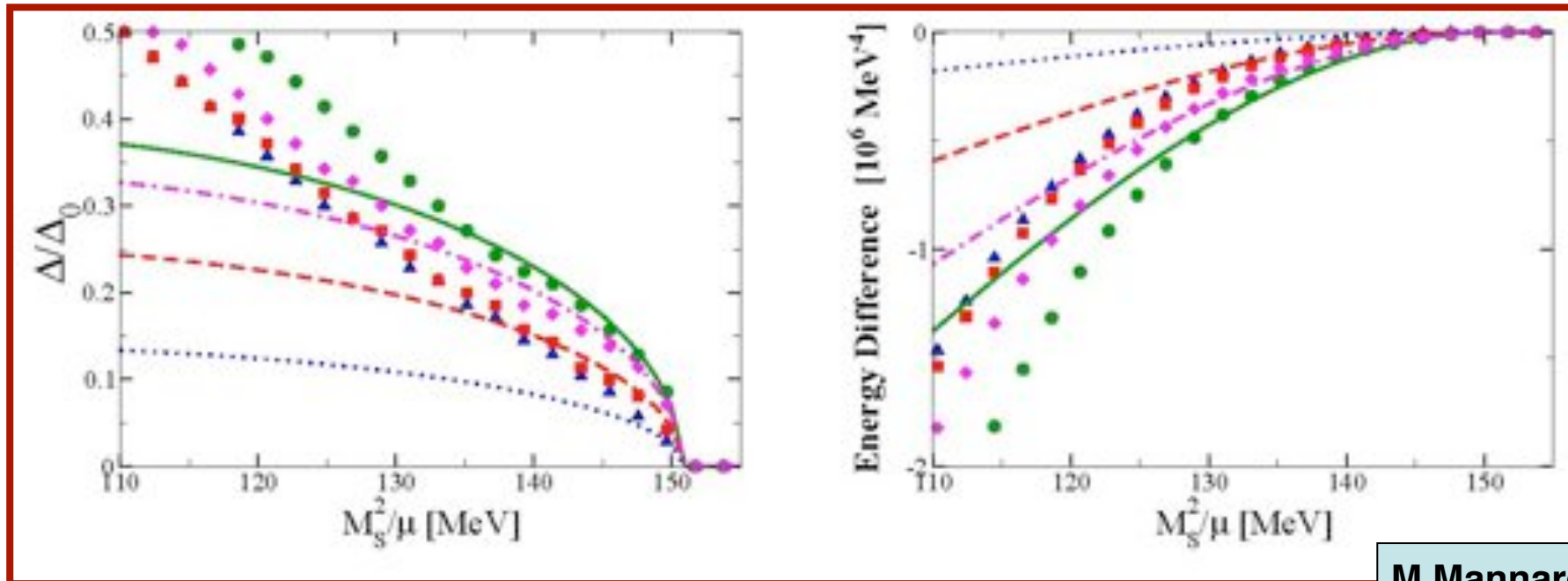
Calculation at leading order in μ



M.Mannarelli,
K. Rajagopal,
R. Sharma

R.Gatto,
R.Casalbuoni,
N.Ippolito,
G.Nardulli,
M.Ruggieri

Comparison between the two approaches



M.Mannarelli,
K. Rajagopal,
R. Sharma

Lines: GL. Dots: complete
Green = $q_2 q_3$ parallel. **Magenta**: $2/3\pi$.
Red; $7/8\pi$. **Blue**: $31/32\pi$

Chromomagnetic instability

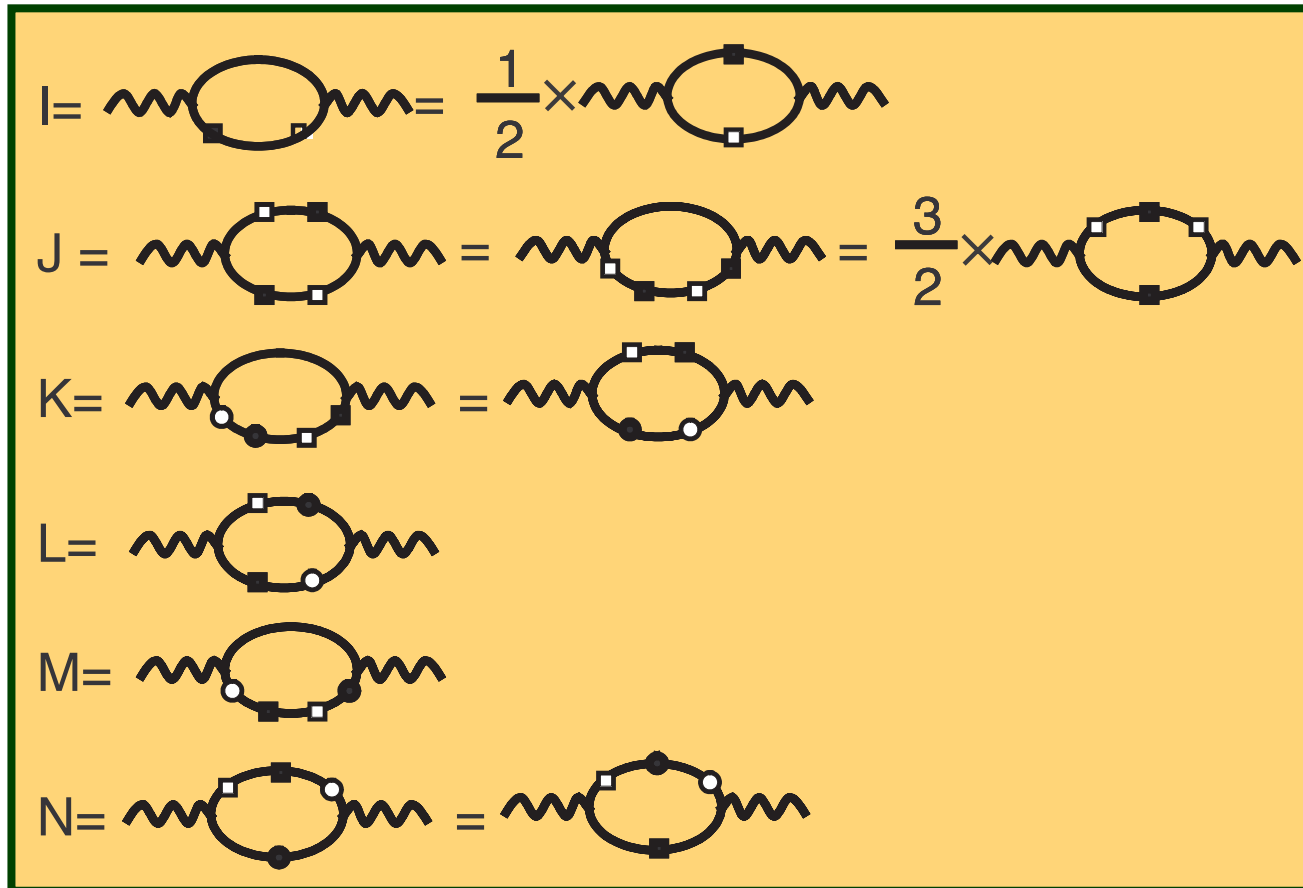
In gCFL : imaginary gluon Meissner masses

Studied in LOFF model with 2 flavors (Giannakis&Ren, Fukushima, Gorbar, Hashimoto&Miransky).

LOFF favored in comparison to gapless uniform conductive states. At least in the GL region (small gap) no chromomagnetic instability

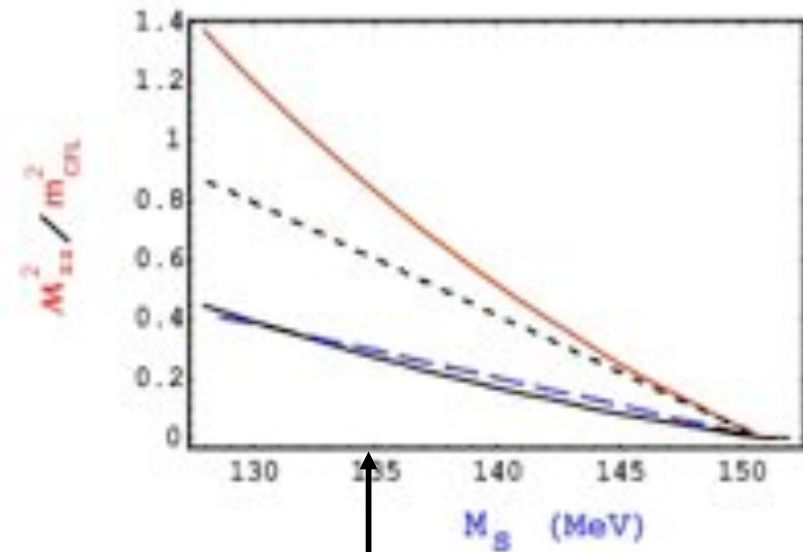
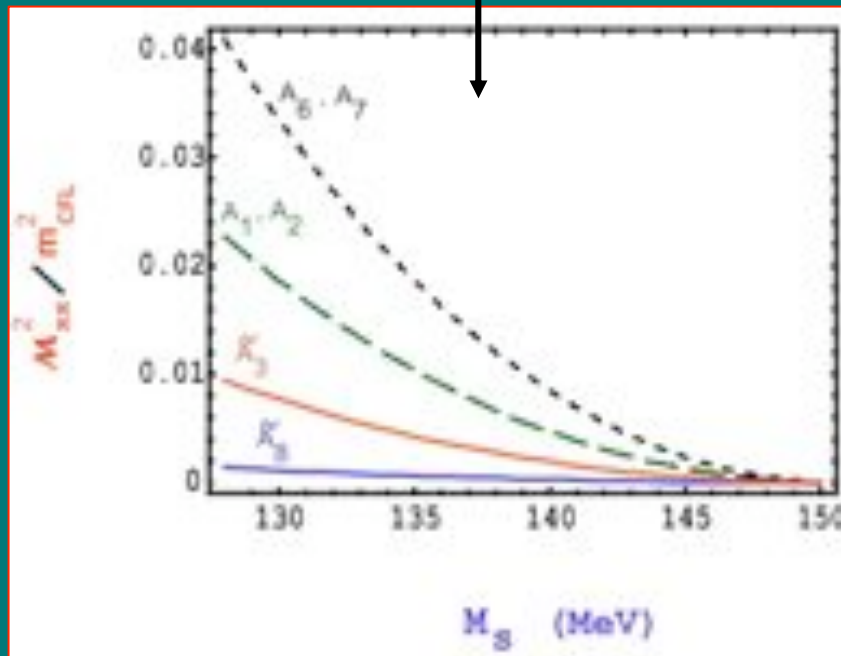
For 3 flavors: recent results (Ciminale, Gatto, GN, Ruggieri)

Feynman Diagrams at $O(\Delta^4)$



Gluon Meissner masses: no instability

Transverse mass



Longitudinal mass

M.Ciminale,
R.Gatto,GN,
M.Ruggieri

Conclusions

- QCD at intermediate densities, small T : active field of research
- Structure of vacuum not yet understood, though color superconductivity expected
- Likely gapless superconductivity, e.g. LOFF anisotropic state
- Results interesting for the effects on cooling of compact stars with a quark core: faster cooling